

Sturm-Liouville Theory (part 2)

$$\frac{d}{dx} \left[p(x) \frac{dy}{dx} \right] + q(x)y + \lambda w(x)y = 0 \quad a \leq x \leq b$$

$$\text{Subject to } \alpha_1 y(a) + \alpha_2 y'(a) = 0$$

$$\beta_1 y(b) + \beta_2 y'(b) = 0$$

the SL problem is said to be regular if on $a \leq x \leq b$,

$p(x)$, $p'(x)$, $q(x)$, $w(x)$ are continuous and

$$p(x) > 0, w(x) > 0$$

→ infinitely many eigenvalues λ that form an increasing sequence

$$\lambda_1 < \lambda_2 < \lambda_3 < \dots < \lambda_n < \dots \quad \lambda_n \rightarrow \infty \text{ as } n \rightarrow \infty$$

(these could be negative but if $\alpha_1, \alpha_2, \beta_1, \beta_2$ are all nonnegative, then λ s are also nonnegative)

→ each λ_n is paired with an eigenfunction y_n and these eigenfunctions are mutually orthogonal with respect to the weight function $w(x)$

$$\int_a^b y_n y_m w(x) dx = 0 \quad \text{if } n \neq m$$

example $y'' + \lambda y = 0$ $0 < x < L$ looks just like

$$y(0) = 0$$

$$\bar{X}'' + \lambda \bar{X} = 0$$

$$y(L) = 0$$

$$\bar{X}(0) = 0$$

$$\bar{X}(L) = 0$$

$$p=1, q=0, w=1$$

$$\alpha_1=1, \alpha_2=0, \beta_1=1, \beta_2=0$$

so, this is a regular SL problem w/ nonnegative λ s

we know $\lambda_n = \frac{n^2 \pi^2}{L^2}$ nonnegative (agrees w/ SL theory)

$$y_n = \sin\left(\frac{n\pi}{L}x\right)$$

$$\int_0^L y_n y_m w(x) dx = \int_0^L \sin\left(\frac{n\pi}{L}x\right) \sin\left(\frac{m\pi}{L}x\right) dx = 0$$

(as predicted by SL theory)

the nonnegative λ is the reason why we assumed the separation constant to be $-\lambda$

$$\frac{\bar{X}''}{\bar{X}} = \frac{T'}{KT} = -\lambda$$

→ Fourier solution is just a special case of SL problem

example

$$y'' + \lambda y = 0$$

$$0 < x < L$$

$$p=1, q=0, w=1$$

$$y(0) = 0$$

$$d_1 = 1, d_2 = 0$$

$$hy(L) - y'(L) = 0 \quad h > 0$$

$$\beta_1 = h > 0, \beta_2 = -1$$

Since not all α, β are nonnegative, the eigenvalues λ s are not always nonnegative

→ negative eigenvalues are possible

let's investigate the (possibly) negative eigenvalues

$$y'' + \lambda y = 0$$

$$\lambda < 0, \text{ for convenience let } \lambda = -k^2 \quad (k > 0)$$

(k^2 avoids $\sqrt{\lambda}$ in solution
convenience only)

$$y'' - k^2 y = 0$$

$$y = C_1 e^{kx} + C_2 e^{-kx}$$

or $y = A \cosh(kx) + B \sinh(kx)$

} choose whichever is more convenient w/ give BCs

$$\text{w/ } y(0)=0 \text{ and } y = c_1 e^{kx} + c_2 e^{-kx}$$

we see $0 = c_1 + c_2$ ok, but it would be nice if one of them is 0

$$\text{w/ } y(0)=0 \text{ and } y = A \cosh(kx) + B \sinh(kx)$$

we see $0 = A \rightarrow y = B \sinh(kx)$ $\lambda = -k^2$
(eigenfunctions are $y_n = \sinh(kx)$)

use $hy(L) - y'(L) = 0$ to find k (and therefore λ)

$$y = B \sinh(kx)$$

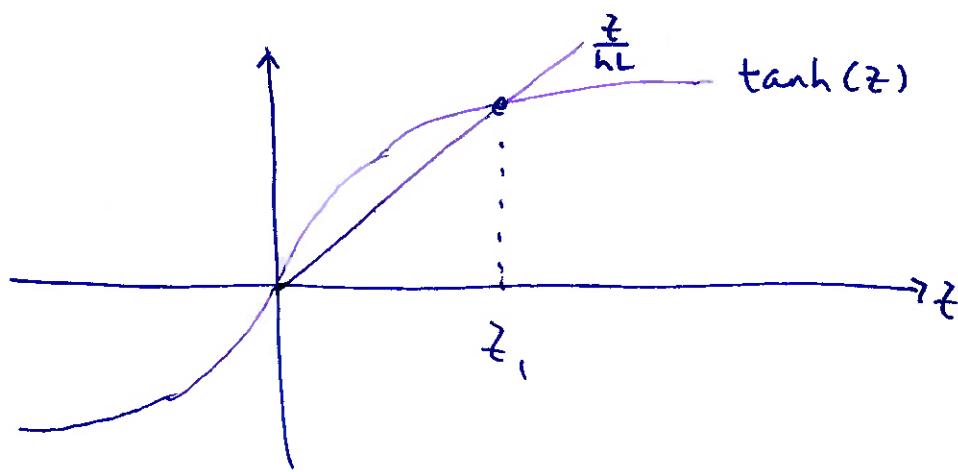
$$y' = BK \cosh(kx)$$

$$hB \sinh(kL) - BK \cosh(kL) = 0 \quad B \neq 0, L \neq 0, h \neq 0$$

$$\tanh(kL) = \frac{k}{h} \quad \text{solve for } k$$

$$\tanh(kL) = \frac{kL}{hL} \rightarrow \text{solve } \tanh(z) = \frac{z}{hL} \quad z = kL$$

intersection of $\tanh(z)$ and
line $\frac{z}{hL}$



at most one intersection
 (if hL is small the
 line is too steep
 to intersect $\tanh(z)$)

at most one intersection \rightarrow at most one negative eigenvalue

$$\lambda_1 < \lambda_2 < \lambda_3 < \dots < \lambda_n < \dots$$

\uparrow negative \uparrow must be ≥ 0

this negative λ is the smallest eigenvalue (no largest λ)

$$\lambda_1 = -k^2 = -\left(\frac{z_1}{L}\right)^2$$

corresponding eigenfunction: $y_1 = \sinh\left(\frac{z_1}{L}x\right)$

next: is $\lambda = 0$ and eigenvalue

then find $\lambda > 0$

let's revisit the heat exchange problem from last time

$$u_t = k u_{xx} \quad 0 < x < L$$

$$u(0, t) = 0$$

$$u_x(L, t) = -h u(L, t) \quad h > 0$$

$u(x, 0) = 100$ initially heated to 100 uniformly

last time we solved the spatial problem: $y'' + \lambda y = 0$

$$\lambda_n = \frac{z_n^2}{L^2}$$

z_n : positive intersections of $\tan(z)$ and $-\frac{z}{hL}$

solve the time problem as usual

then the general solution

$$u(x, t) = \sum_{n=1}^{\infty} c_n e^{-k\lambda_n t} \sin(\sqrt{\lambda_n} x)$$

$$u(x, 0) = 100$$

$$100 = \sum_{n=1}^{\infty} C_n \sin(\sqrt{\lambda_n} x)$$

NOT a sine series since

λ is not integer multiple of $\frac{\pi}{L}$

$$\text{so, } C_n \neq \frac{2}{L} \int_0^L 100 \sin\left(\frac{n\pi}{L} x\right) dx$$

we can still use the orthogonality of eigenfunctions (guaranteed by SL theory) to find C_n

$$\int_0^L y_n y_m w(x) dx = 0 \quad n \neq m$$

here, $w(x) = 1$

$$y_n = \sin(\sqrt{\lambda_n} x)$$

find C_n :

$$100 \sin(\sqrt{\lambda_m} x) = \sum_{n=1}^{\infty} C_n \sin(\sqrt{\lambda_n} x) \sin(\sqrt{\lambda_m} x)$$

integrate over $0 < x < L$

$$\int_0^L 100 \sin(\sqrt{\lambda_m} x) dx = \sum_{n=1}^{\infty} \underbrace{\int_0^L C_n \sin(\sqrt{\lambda_n} x) \sin(\sqrt{\lambda_m} x) dx}_{0 \text{ except } n=m}$$

$$\int_0^L 100 \sin(\sqrt{\lambda_n} x) dx = \int_0^L C_n \sin^2(\sqrt{\lambda_n} x) dx$$

$$C_n = \frac{\int_0^L 100 \sin(\sqrt{\lambda_n} x) dx}{\int_0^L \sin^2(\sqrt{\lambda_n} x) dx}$$

if $\lambda_n = \frac{n^2 \pi^2}{L^2}$ it reduces to $\frac{2}{L} \int_0^L 100 \sin\left(\frac{n\pi}{L} x\right) dx$